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## PHASE-DEPENDENT SPECTROSCOPY OF MIRA VARIABLE STARS

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### ABSTRACT

Spectroscopic measurements of Mira variable stars as a function of phase probe the stellar atmospheres and underlying pulsation mechanisms. For example, measuring variations in TiO, VO, and ZrO with phase can be used to help determine whether these molecular species are produced in an extended region above the layers where Balmer line emission occurs or below this *shocked* region. Using the same methods, the Balmer line *increment*, where the strongest Balmer line at phase zero is H $\delta$  and not H $\alpha$ , can be measured and explanations tested, along with another peculiarity, the absence of the H $\epsilon$  line in the spectra of Mira variables when the other Balmer lines are strong. We present new spectra covering the spectral range from 6200 to 9000 Å of 20 Mira variables. A relationship between variations in the Ca II IR triplet and H $\alpha$  as a function of phase support the hypothesis that H $\epsilon$ 's observational characteristics result from an interaction of H $\epsilon$  photons with the Ca II H line. New periods and epochs of variability are also presented for each star.

*Key words:* Miras—techniques: spectroscopic

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### 1. INTRODUCTION

Mira-type variable stars are large, cool stars whose visual light variations exceed 2.5 mag over periods from 150 to  $\sim$ 500 days. The light curves of Mira variables depend on the surface temperature, radius, and opacity, all of which vary as the star pulsates. These pulsations extend the atmosphere beyond that of the hydrostatic equilibrium configuration and enhance mass loss in these stars ([Maciel 1977](#); [Willson & Hill 1979](#); [Bertschinger & Chevalier 1985](#); [Bowen 1988](#); [Fleischer et](#)

al. 1992, 1995). As a result, Mira variables are an important component in seeding the interstellar medium with C, N, and O.

These stars are located on the asymptotic giant branch, a transitional phase in stellar evolution. Photometric and spectroscopic measurements of their light curves provide a means to probe the stellar atmospheres and underlying pulsation mechanisms occurring during this stellar phase. In the near infrared, the spectra of Mira stars are dominated by the TiO  $\gamma$  system, the VO  $\gamma$  system, and ZrO (Wing 1967). The TiO features are thought to be produced in a layer somewhat far from the photosphere (Gillet 1988). Haniff et al. (1992) present optical aperture synthetic images of the photosphere of Ceti at 6500 Å, 7007 Å, and within a TiO band head at 7099 Å, with the star phase  $\sim 0.94$ . They find asymmetry in the images, with the TiO image  $1 \frac{1}{2}$  times larger than the photospheric images. Also, narrowband speckle interferometric measurements taken in the TiO 7120 Å band head and outside at 7400 Å by Labeyrie et al. (1977) show that the diameters of R Leo and Cet are twice as large in the TiO feature than outside of it. This demonstrates that a model atmosphere, based on the spectra observed over a TiO bandpass, provides parameters such as  $T_{\text{eff}}$  and  $\log g$  in an atmospheric layer far from the photosphere.

Joy (1926) presents a comprehensive, phase-dependent, spectroscopic ( $35 \text{ \AA mm}^{-1}$ ) study of a prototype Mira variable,  $\bar{\nu}$  Ceti. Analyzing 131 spectra taken over a 10 yr period, Joy describes several important characteristics. Briefly, the spectra of Ceti show that TiO bands vary with magnitude, hydrogen emission lines appear with greatest intensity at or shortly after maximum visual brightness (phase zero), ionized iron emission lines are observed at maximum, and the temperature is estimated to vary from 1800 to 2300 K. Absorption lines (including iron, vanadium, chromium, manganese, calcium, and magnesium) were used to measure a variation in radial velocity. The maximum positive velocity occurs at phase zero, and the greatest blueshift occurs at minimum light. Later, Joy (1954) took 88 spectrograms of Ceti at a higher spectral resolution (typically  $10.3 \text{ \AA mm}^{-1}$ ) and confirmed that maximum velocity of recession occurs soon after visual maximum. He attributed these results to a pulsational mechanism and suggested the possibility of shocks.

Perhaps one of the most interesting characteristics of Mira spectra is the strong hydrogen Balmer line emission that is seen throughout much of the pulsation cycle. As Pickering (1887) first noticed in spectra of Ceti and as was later described in detail by Joy (1926, 1947, 1954), the hydrogen Balmer emission-line flux in relation to the nearby *photospheric* (i.e., pseudocontinuum) flux is unique in the oxygen-rich (M-type) Mira spectra: Balmer  $H\alpha$  emission is typically weaker than  $H\beta$ , which, in turn, is weaker than  $H\gamma$  near peak visual brightness.  $H\delta$  is seen as the strongest Balmer emission line at phase zero. Lines higher in the series (i.e., toward shorter wavelengths) are weaker. This *Balmer increment* (i.e.,  $F_{H\alpha} < F_{H\beta} < F_{H\gamma} < F_{H\delta}$ ) is just opposite of what would be expected.  $H\alpha$ , having the largest oscillator strength, should be stronger than  $H\beta$ , and the higher order Balmer lines should be weaker down the line (i.e., one should see a Balmer line *decrement*), assuming these lines all form in the same region of the atmosphere (i.e., similar  $T$  and  $P$ ). Meanwhile, in S-type and carbon star (N-type) Mira variables, the strength of the Balmer lines approximately follow their expected respective oscillator strengths (Merrill 1940, p. 44).

For years, this Balmer line increment in the M-type Mira variables has been attributed to TiO

absorption, which hides  $H\alpha$ ,  $H\beta$ , and  $H\gamma$  fluxes ([Merrill 1940](#); [Joy 1947](#); [Gillet 1988](#)), although there is some debate of the effect that this or other molecular absorption has on the  $H\alpha$  line ([Gillet, Maurice, & Baade 1983](#); [Gillet et al. 1985](#)). Recently, another explanation has been given for this Balmer increment: non-LTE (NLTE) radiative transfer calculations of hydrodynamic models representative of Mira variables ([Bowen 1988](#)) suggest that the Balmer increment results from radiative transfer effects in the hydrogen lines themselves when formed in a shocked atmosphere ([Luttermoser & Bowen 1992](#); [Luttermoser, Bowen, & Willson 2000](#)). In these calculations,  $H\alpha$ , having the highest optical depth, forms just in front of the innermost shock.  $H\beta$  then forms a little deeper as a result of its lower optical depth, and  $H\gamma$  deeper still. The optical depth of  $H\delta$  causes it to arise from the hottest part of the shock, and hence it extends higher above the continuum than the longward Balmer lines. Then, as one goes to higher order lines in the series, the opacity in these lines is not high enough in the shock for these lines to form there—we see through the shock at these transitions. As such, this increment may be giving us information on the shock thickness for Mira variables and may indicate that the shock structure of the S-type and N-type Mira variables is fundamentally different from that of the M-type Mira variables. Future NLTE radiative hydrodynamic models of Mira variables with different C/O ratios are needed to see if this is the case. It is likely that a combination of both processes (i.e., TiO absorption and NLTE radiative transfer effects) are responsible for this Balmer line increment.

Another striking feature is the weakness (and often absence) of the  $H\epsilon$  line (3970.074 Å) in the Balmer series near maximum visual brightness (see [Gillet 1988](#)). [Merrill \(1940\)](#) noted this and suggested this weakness was a result of the interaction between the  $H\epsilon$  transition and the Ca II H line (3968.470 Å) wing. [Castelaz & Luttermoser \(1997, hereafter CL97\)](#) concur with this suggestion. Briefly, the  $H\epsilon$  photons may be scattered by the Ca II H line out to IR wavelengths via the 8662 Å line. Of the three lines in the Ca II IR triplet, the 8498 and 8542 Å lines share the same upper level of the Ca II K line, whereas the 8662 Å line shares the same upper level as the Ca II H line. As such, if  $H\epsilon$  photons are being scattered by Ca II H, the  $\lambda 8662$  will show variations independent of the other two Ca II IR triplet lines and can be tested by monitoring the strength of the absorption of this 8662 Å line as compared with the 8498 and 8542 Å lines as a function of phase in the Mira stars.

Ca II IR triplet observations of Mira variables have been carried out by [Contadakis & Solf \(1981\)](#) and [Gillet et al. \(1985\)](#). Unfortunately, [Contadakis & Solf \(1981\)](#) only made one observation of the 8662 Å line in their monitoring program of S-type Mira variables. They observed  $\lambda 8498$  and  $\lambda 8542$  in emission in many of the stars in their sample near phase zero and as absorption lines at other phases, so we cannot use their data to test our proposed hypothesis. [Gillet et al. \(1985\)](#) observed P Cygni profiles in the Ca II IR triplet lines near phase zero in the hot Mira variable S Car. Once again, only one spectrum of the 8662 Å line was obtained in their sample of Ca II spectra. [CL97](#) present a set of phase-dependent spectra taken especially to address this problem of  $H\epsilon$  photon scattering by Ca II. Spectra of seven Mira variables were taken at different phases and suggest a possible anticorrelation between  $H\alpha$  emission and Ca II  $\lambda 8662$  absorption. Assuming that the  $H\epsilon$  line-strength variations are in phase with  $H\alpha$ , then the apparent anticorrelation between the strength of the  $H\alpha$  emission line and the strength of the Ca II 8662 Å line suggests that a *fluorescence* is taking place in the Ca II 8662 Å line with  $H\epsilon$  serving as the pump through the Ca II H line. In this paper we present phase-dependent spectra of 20 Mira variables to further explore the anticorrelation of Ca II  $\lambda 8662$  absorption with  $H\alpha$  emission.

## 2. OBSERVATIONS

The 20 stars for which spectra are presented, their equatorial coordinates, mean spectral types, and visual maxima and minima taken from the SIMBAD database are given in [Table 1](#). Also listed in [Table 1](#) are new ephemerides of these Mira stars. We calculated the ephemerides from AAVSO light curves measured within the previous 8 yr. The curves were fitted with a linear combination of sine and cosine functions from which a new period and Julian date at phase zero were determined for each star.

TABLE 1 MIRA VARIABLES

Spectra of Mira variable stars were taken between 1997 February and 1998 November using a low-resolution spectrograph. Since we are interested in only monitoring absorption- and emission-line strengths for this program, low-resolution spectra are all that is required. The spectrograph was used at both the Southeastern Association for Research in Astronomy (SARA) 0.9 m telescope at Kitt Peak, and Appalachian State University's Dark Sky Observatory (DSO) 0.45 m telescope located near Boone, North Carolina. A converter lens was used at both sites to convert the respective telescope f-ratios to about f/11 for the spectrograph.

The spectrograph was configured with a 600 groove  $\text{mm}^{-1}$  grating. The slit width was 100  $\mu\text{m}$ . At the SARA 0.9 m telescope, the slit width was 3'', and at the DSO 0.45 m telescope the slit width was 4''. At both sites, the slit was parallel to the hour angle. A cooled 768  $\times$  512 CCD camera with 9  $\mu\text{m}$  square pixels was used to record the spectra. The linear dispersion was 1.08  $\text{\AA pixel}^{-1}$ , covering 768 pixels or 829  $\text{\AA}$  on the CCD. The spectral resolution was measured to be 2.4  $\text{\AA}$ . By rotating the grating up to four times per star, spectra were taken from about 6200 to 9000  $\text{\AA}$  and included H $\alpha$ , TiO, VO, and the Ca II infrared triplet lines. Integration times were adjusted to achieve a signal-to-noise ratio greater than 100 for most spectra, except for the 1997 May spectra of W Her, T Her, and W Lyr, and the 1998 May spectrum of T Her.

[Table 2](#) gives the log of observations for each star, which includes dates of observation, phase of the variable, approximate visual magnitude, integration time, and observing site. The phases listed in [Table 2](#) were determined from the ephemerides given in [Table 1](#) and refer to the visual phases, with phase zero corresponding to maximum visual brightness. The visual magnitudes have been obtained from curve fits to the light curves from the AAVSO database. The method of phase determination used in this paper is not the standard method used by AAVSO. We used the AAVSO light curves observed at the same time as our observations to determine the phases. The AAVSO standard method to derive official dates of maximum and minimum may give different phases.

TABLE 2 LOG OF OBSERVATIONS

Dark frames and flat frames were taken for flat-fielding purposes. Spectra of neon lamp emission were taken simultaneously with the stellar spectra and used for wavelength calibration. We flat-fielded the images and extracted the spectra using MIRA software. The extracted spectra were wavelength-calibrated using the spectrum of neon superposed on the CCD frame with the stellar spectrum (Crowe, Heaton, & Castelaz 1996).

### 3. RESULTS

The spectra for the Mira variable stars are shown in Figure 1. The wavelengths of TiO, VO, ZrO, the Ca II IR triplet, H $\alpha$ , and terrestrial oxygen are marked above each set of spectra. Their wavelengths were identified in CL97. The ZrO and VO absorption overlap at 6574 and 6578 Å, and VO and TiO absorb at 7865 and 7861 Å, respectively. Because of the low dispersion of our spectra, these features are blended. The appearance of a relatively narrow feature near 8230 Å is seen in some of the spectra—an  $\bar{\nu}$  Cet spectrum taken 1997 November 1 (phase 0.80), and spectra taken 1997 May 26–28 of R Leo (0.49), V CVn (0.65), R Boo (0.11), R Ser (0.26), RS Her (0.75), and W Lyr (0.94). This feature is due to terrestrial H<sub>2</sub>O at 8227 Å (Turnshek et al. 1985) and is an effective measure of the relative humidity in the air.

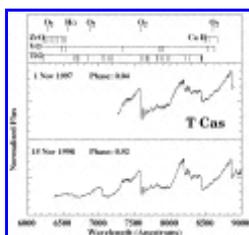
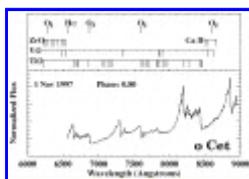
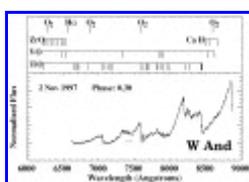
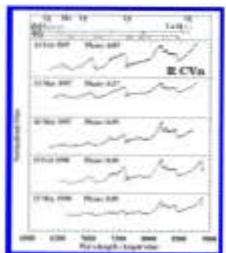
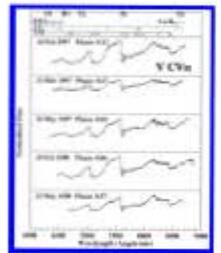
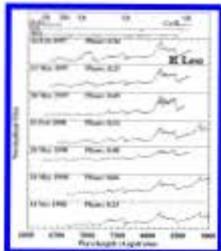
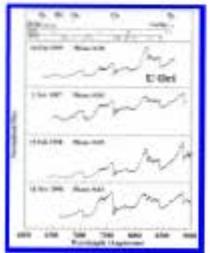
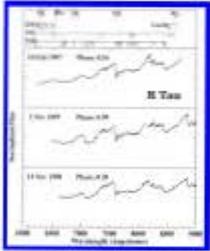
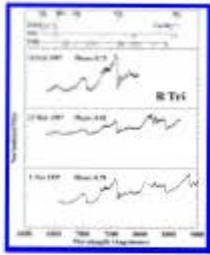
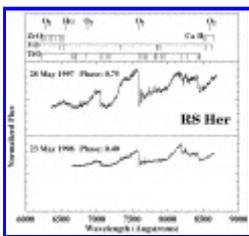
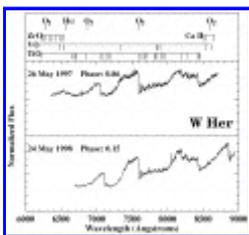
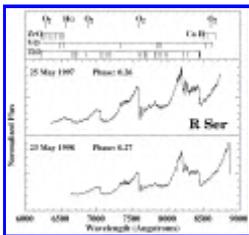
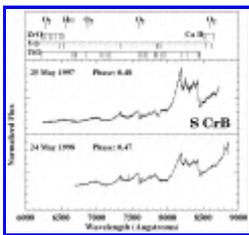
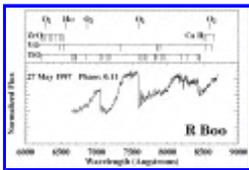
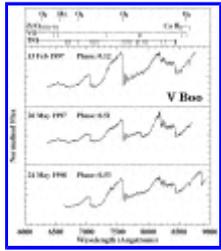
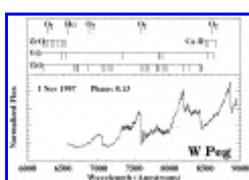
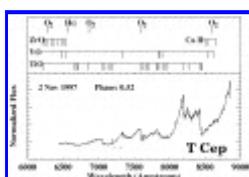
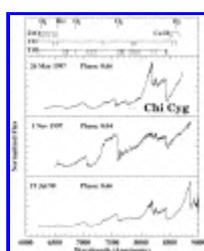
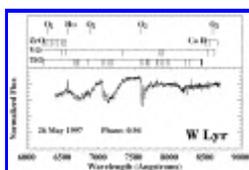
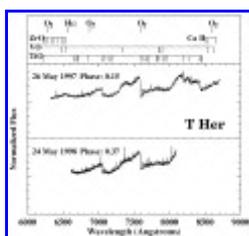


FIG. 1.—Spectra of the 20 Mira variables. The name of the star, date, and phase are given on each spectrum. The flux is normalized to 1. Above each set of spectra are markers for the major spectral features, as well as terrestrial oxygen. H $\alpha$  is weak or not seen in most of the spectra. Markers enclose regions of individual spectra where no data were taken. Spikes due to cosmic-ray hits in several spectra have not been removed.









### 3.1. Radiative Transfer in the Ca II Ion

The transfer of radiation in the Ca II ion is very complicated. Besides H $\epsilon$  photons affecting the level populations of the  $3d\ ^2P_{1/2}$  state through the Ca II H line, the hydrogen Ly $\alpha$  line lies just shortward of the Ca II ionization edge of the  $3d\ ^2D_{3/2}$  (at 1218.1 Å) and the  $3d\ ^2D_{5/2}$  (at 1219.0 Å) states. These two states are the lower levels of the Ca II IR triplet lines and are metastable. If Ly $\alpha$  is a strong emission feature, and the Ca II  $^2D$  continuous opacity forms in a region of the atmosphere where the Ly $\alpha$  line is not yet in detailed balance, then Ly $\alpha$  photons may influence the Ca II IR triplet lines as well. Note that no observations have yet been made in the far-UV for Mira variables.

The line center of Ly $\alpha$  lies 2.4 Å shortward of the  $^2D_{3/2}$  ionization edge and 3.3 Å shortward of the  $^2D_{5/2}$  edge, so much of the Ly $\alpha$  emission profile can ionize Ca II out of the metastable state. The

lower level of the Ca II 8662 Å line transition is the  $^2D_{3/2}$  state, whose ionization edge lies slightly closer to Ly $\alpha$  than the  $^2D_{5/2}$  edge. The following question arises: Will Ly $\alpha$  photons affect the level densities of the two  $3d^2D$  states differently? To answer this question, we ran a few atmospheric models with an arbitrarily located 10,000 K shock through the LTE stellar atmosphere code ATLAS (Kurucz 1970; Brown et al. 1989). Although it is likely that NLTE effects will dominate the level and ion densities in the atomic and molecular species (Luttermoser & Bowen 1992; Luttermoser, Bowen, & Willson 2000), these LTE runs are performed to merely determine the variation of the Ca II  $^2D$  continuous opacity across the Ly $\alpha$  profile. It also should be pointed out that in regions of the atmosphere where LTE no longer applies, the assumptions of radiative and hydrostatic equilibria are no longer valid in these pulsating giant stars either. In fact, a Mira star has numerous shocks propagating through its atmosphere at any given time, as has been shown by Bowen (1988) and more recently by Höfner et al. (1998) and Loidl et al. (1999). Willson (2000) gives a very detailed review of all the dynamic modeling that has been performed on these pulsating stars and discusses the problems of carrying out NLTE radiative transfer in such a dynamic atmosphere.

We sampled atmospheric depths in front of the shock, in the shock, behind the shock, and deep in the photosphere where the continuum reaches optical depth unity in this region of the spectrum. We found that the continuous opacity from the Ca II  $3d^2D$  ionization remained constant (from both  $J$  sublevels) to within 0.15% from the location of the edges through 1210 Å, which should include most of the Ly $\alpha$  emission profile. The fact that the continuous opacity from Ca II remains constant across the Ly $\alpha$  profile indicates that Ly $\alpha$  will not preferentially affect the number density in the  $^2D_{3/2}$  level as compared with the  $^2D_{5/2}$  level—photoionizations of Ca II due to Ly $\alpha$  photons will affect the strengths of three Ca II IR lines in a similar fashion. Therefore, any variation in the 8662 Å line that is not seen in the other lines must result from some process other than Ly $\alpha$  photoionizations.

### 3.2. The H $\alpha$ Emission Line and Ca II IR Triplet

We are interested in the strength of the Ca II IR triplet compared with the strength of H $\alpha$  as a function of phase, since we are assuming that variations in H $\epsilon$  will mimic variations in H $\alpha$ . The Ca II IR triplet lines are not strong, as expected for stars later than M0 (Zhou 1991). A total of 27 spectra of 15 stars in our sample span the wavelength range from H $\alpha$  through the Ca II IR triplet. The remaining spectra are missing either the H $\alpha$  or the Ca II IR triplet regions of the spectrum because the spectrograph grating was not rotated sufficiently during observation to cover those parts of the spectrum.

The observations that show obvious H $\alpha$  emission features include U Ori (phase 0.29), R Leo (0.34), V CVn (0.12), R CVn (0.09 and 0.17), R Ser (0.26), W Lyr (0.94), and X Cyg (0.04). At the same phases, the Ca II  $\lambda$ 8662 is seen in emission in R Leo, R Ser, W Lyr, and X Cyg, whereas the 8498 Å and 8542 Å lines stay in absorption or are not apparent. Ca II  $\lambda$ 8662 is not seen in emission in any other spectra that we took, only in those that show H $\alpha$  emission. The Ca II IR triplet can be seen in absorption in the remaining 23 spectra (although in some cases weakly), except for U Ori (phase 0.01), S CrB (0.48), W Her (0.86), and T Her (0.15), which do not appear to have any type of Ca II IR triplet features.

At this point, we call attention to three of the hydrogen Paschen lines that lie close to each Ca II IR

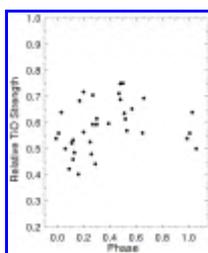
triplet line:  $\lambda 8502.4$  (4.4 Å redward from the Ca II 8498.0 Å line),  $\lambda 8545.3$  (2.8 Å redward of the Ca II 8542.1 Å line), and  $\lambda 8665.0$  (2.8 Å from the Ca II 8662.1 Å line), the Pa13, Pa12, and Pa10 lines, respectively. It has been shown by [Gillet et al. \(1985\)](#) that even though the Pa $\delta$  line is in emission in the spectrum of the Mira variable S Car, the higher order Paschen lines near the Ca II IR triplet are seen neither in absorption nor in emission, analogous to the Balmer lines. As such, it is unlikely that these higher order Paschen lines are affecting the Ca II lines in our spectra.

### 3.3. The TiO Bands and H $\alpha$ Emission

In addition to the Ca II IR triplet, we are interested in comparing the TiO  $\gamma$  system's molecular features with H $\alpha$  emission as a function of phase. A qualitative comparison of H $\alpha$  emission with the molecular features can be made for U Ori, R Leo, R CVn, R Ser, W Lyr, and  $\chi$  Cyg, stars with phase-dependent spectra that also show the H $\alpha$  emission line in at least one spectrum. Weak H $\alpha$  emission is seen in U Ori at phase 0.29 on 1998 February 15. However, on the previous pulsation cycle, H $\alpha$  emission is not obvious at either phases 0.30 or 0.01, although there may be some very weak emission. Stronger H $\alpha$  emission is seen in R Leo (phase 0.34), R CVn (0.09, 0.17, and 0.39), R Ser (0.26), W Lyr (0.94), and  $\chi$  Cyg (0.04). As mentioned above, H $\alpha$  is notorious for being observed as a weak emission feature when H $\beta$ , H $\gamma$ , and H $\delta$  are strong. Surprisingly though, H $\alpha$  emission was not seen in various spectra where we would expect to find it: R Tau (phase 0.16 and 0.99), R Leo (0.16 and 0.25), V CVn (0.12, 0.27, and 0.06), R CVn (0.20), V Boo (0.12), and T Her (0.15).

[Merrill \(1940\)](#) and [Joy \(1926\)](#) report that the TiO bands are regularly stronger at minimum than at maximum light in Mira variables. To measure this trend, we checked the variability of the TiO feature at 7054 Å (band head) with respect to a portion of the flux uncompromised by TiO, VO, and ZrO. We integrated the flux of each spectrum in the 6995–7045 Å wavelength band (non-TiO) and the 7060–7110 Å band (TiO), each 50 Å wide. We then divided the integrated flux of the TiO band by the non-TiO band. By doing this, any scattered light that may have been in the spectra is effectively canceled out. Since our data set contains warm oxygen-rich (earlier than M6), cool oxygen-rich (M6–M8), and one MS Mira sampled sporadically over various phases, we only include the cool oxygen-rich Mira variables in [Figure 2](#), which graphs the aforementioned flux ratio as a function of phase. The M6–M8 stars are selected here to minimize the spread of effective temperatures at maximum light, which will influence the strength of the TiO bands. Ideally, one would want many observations of each star over a single pulsation cycle. However, with the limited sample we have, we feel that we can obtain an approximate test of TiO variation with respect to phase. There appears to be no obvious trend in variations in the TiO 7054 Å band flux with respect to the non-TiO band flux. Because of this observation, variations seen in H $\alpha$  as a function of phase must primarily result from variations in the H $\alpha$  emission itself and not from variations in overlying TiO absorption. As a result, using H $\alpha$  flux variations as a proxy for variations in the intrinsic H $\epsilon$  flux is valid from this analysis.

FIG. 2.—Integrated flux ratio of the 7060–7110 Å wavelength band (A) to that of the 6995–7045 Å band (B) plotted as a function of phase. Band A contains TiO opacity whereas band B is free of TiO, VO, and ZrO band heads. This plot contains only data from the M6, M7, and M8 spectral type Mira variables. No apparent trends are seen in the data, which indicates that variations in the H $\alpha$  flux result primarily from intrinsic flux variations in H $\alpha$  and not from varying overlying TiO absorption.



#### 4. DISCUSSION

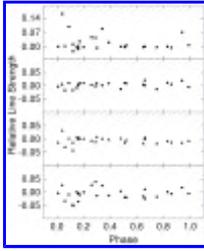
We wish to test the idea that the apparent lack of H $\epsilon$  emission at 3970 Å, when the other Balmer lines are strong emission features, is anticorrelated with the strength of the Ca II absorption line at 8662 Å. As reported in § 1, this anticorrelation results from H $\epsilon$  photons being scattered by the Ca II H line out to the Ca II line at 8662 Å, causing this Ca II absorption line to be *filled in* with respect to the other two Ca II IR triplet lines.

We use H $\alpha$  as a proxy for the H $\epsilon$  line. Following the analysis presented by CL97, we determined a relative *line strength*,  $F$ , for H $\alpha$  emission and the Ca II IR triplet absorption lines. Two points were selected on either side of the emission or absorption feature. The wavelengths of these points were kept constant for all measurements. The observed *profile* was integrated across the wavelength window defined by these two points, resulting in an integrated flux  $f_i$ . A straight line connected between these two points represents a *pseudocontinuum*, and the integrated flux,  $f_c$ , is calculated for it. Then

$$F = (f_i - f_c)/f_c,$$

where  $F$  will be negative for absorption lines and positive for emission lines. The measurements were done for the stars in our sample, where their spectra included wavelengths below 6563 Å and above 8662 Å; a total of 27 spectra. Figure 3 shows the relative line strengths of H $\alpha$  and the Ca II 8498, 8542, and 8662 Å absorption lines as a function of phase. The uncertainty in the measurements is about  $\pm 0.007$ . The relative line strength of H $\alpha$  clearly shows a large scatter near visible maximum and is zero within the uncertainty from phase 0.5 to 0.8. This is consistent with Balmer emission lines becoming prominent near maximum visible light. The Ca II 8498 Å line strengths are zero, within the uncertainty of the measurements; variation is not observed in this line. The Ca II 8542 Å line does show some scatter near phase 0.0, and is zero after phase 0.2. The Ca II 8662 Å line shows more scatter than either of the other two Ca II IR triplet lines, particularly near phase 0.0.

FIG. 3.—Relative line strengths of H $\alpha$  (*squares*) and Ca II lines at 8498 Å (*circles*), 8542 Å (*triangles*), and 8662 Å (*diamonds*) as a function of light variation phase. Variation in the Ca II lines at 8498 and 8542 Å mimic each other, whereas the Ca II line at 8662 Å does not follow the same trend. The uncertainty of the measurements is  $\pm 0.007$ .



From [Figure 3](#), it is difficult to see any correlation between the relative line strength of the  $H\alpha$  and Ca II IR triplet lines. However, we can plot the relative line strengths of Ca II IR triplet lines versus the strength of the  $H\alpha$  line to look for correlations. [Figure 4](#) shows the relative line strengths of Ca II  $\lambda 8498$ ,  $\lambda 8542$ , and  $\lambda 8662$  versus the relative line strength of  $H\alpha$ . The data of each plot are linearly fitted and the results of the fits are drawn in the plots. The slopes of the Ca II  $\lambda 8498$  and  $\lambda 8542$  relative line strengths versus  $H\alpha$  are 0.08 and 0.05, respectively. The slope of the linear fit of the Ca II  $\lambda 8662$  relative line strength versus  $H\alpha$  relative line strength, on the other hand, is 0.32, which is significantly different than the other two Ca II IR triplet linear fits. Furthermore, the Ca II  $\lambda 8662$  versus  $H\alpha$  relative line strength slope is positive, so that as the strength of the  $H\alpha$  line increases, the strength of the Ca II 8662 Å line decreases (i.e., becomes more positive) and even goes into emission. This is the effect we expect if He photons are being scattered by the Ca II H line out to the Ca II line at 8662 Å, causing the Ca II 8662 Å absorption line to be filled in with respect to the other two Ca II IR triplet lines.

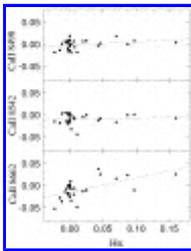


FIG. 4.—Relative line strengths of Ca II  $\lambda 8498$ ,  $\lambda 8542$ , and  $\lambda 8662$  vs. the relative line strength of  $H\alpha$ . The data of each plot are linearly fitted, and the results of the fits are drawn in the plots. Only the Ca II  $\lambda 8662$  vs.  $H\alpha$  plot shows a significant slope, which implies a correlation between the occurrence of the two features.

## 5. CONCLUSION

The 6200 to 9000 Å spectra of Mira variables taken at different phases support a possible anticorrelation between  $H\alpha$  emission and Ca II  $\lambda 8662$  absorption as first suggested by [CL97](#). Assuming that the He line strength variations are in phase with  $H\alpha$ , then the apparent anticorrelation between the strength of the  $H\alpha$  emission line and the strength of the Ca II 8662 Å line suggests that a *fluorescence* is taking place in the Ca II 8662 Å line, with He serving as the pump through the Ca II H line. This type of fluorescence is common in Mira type variables. The strong Fe I (42) lines at 4202 and 4308 Å seen in Mira variables are well-known fluoresced features; in this case, the ultraviolet Mg II *h* and *k* lines serve as the pump via an Fe I (UV3) transition (e.g., [Bidelman & Herbig 1958](#); [Willson 1976](#); [Luttermoser 1996](#)).

The next phase of this research program is to systematically determine  $T_{\text{eff}}$  and  $\log g$  as a function of phase for the Mira variables in our sample. A detailed description of the LTE model synthetic spectra is given by [Piontek & Luttermoser \(2000\)](#).

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*Note added in proof.*—The method of phase determination used in this paper is not the standard method used by the AAVSO. We used the AAVSO light curves observed at the same time as our observations to determine the phases. The AAVSO standard method to derive official dates of maximum and minimum may give different phases.

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